

# 1 Evidence for an Explosive Origin of Central Pit Craters on Mars

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4

## 5 **Abstract:**

6 Kilometer-scale pit craters are nested in the centers of many impact craters on Mars as  
7 well as on icy satellites. They have been inferred to form in the presence of a water-ice rich  
8 substrate; however, the process(es) responsible for their formation is still debated. Leading  
9 models invoke origins by either explosive excavation, or by subsurface drainage and collapse. If  
10 explosive excavation forms central pits, ejecta blankets should be draped around the pits,  
11 whereas internal collapse should not deposit significant material outside pit rims. Using visible  
12 wavelength images from the MRO CTX and HiRISE instruments and thermal infrared images  
13 from the Odyssey THEMIS instrument, we conducted a survey to characterize, in detail, the  
14 global population of central pits in impact craters  $\geq 10$  km in diameter. We specifically examined  
15 the morphology and thermophysical characteristics of the pits for evidence of pit ejecta. Our  
16 analysis of thermal images suggests that coarse-grained materials – which we interpret as pit  
17 ejecta – are distributed proximally around many central pits on the floors of their parent craters.  
18 These observations and interpretations support an explosive origin for central pits on Mars. We  
19 present an alternative “uplift contact model” to explain the formation of central pits late in the  
20 impact process. Theoretical calculations show that more than enough thermal energy is available  
21 via impact melt from the parent crater to form central pits by steam explosions, and such  
22 explosions would require only a modest amount (2-6% by volume) of uplifted water-ice. We

23 therefore propose that central pits on Mars could have formed explosively by the interaction of  
24 impact melt and subsurface water-ice.

25

## 26 **Introduction:**

27         Nested central pit craters on Mars have remained enigmatic structures for several  
28 decades. Their formation is typically believed to be connected with subsurface water due to their  
29 relative abundance on Mars and icy satellites, but the exact role of water and the specific  
30 process(es) responsible for forming central pits are still debated. In this study, we make thermal  
31 inertia observations of central pit craters to test hypotheses for central pit formation. We start  
32 with an overview of the previously proposed hypotheses and the gaps in our understanding.  
33 Then, we discuss the utility of thermal inertia in remotely determining grain size distributions  
34 around central pits. We hypothesize that central pits are formed by explosive excavation or  
35 devolatilization during or after impact. After analyzing our results, we present an alternative  
36 model for central pit formation that -- uniquely among other pit origin hypotheses -- creates an  
37 explosion late enough in the impact process for central pits to be preserved. Finally, we apply our  
38 integrated observations to interpret the morphology and thermal properties of central pits in the  
39 context of central uplifts and propose testable predictions for the model.

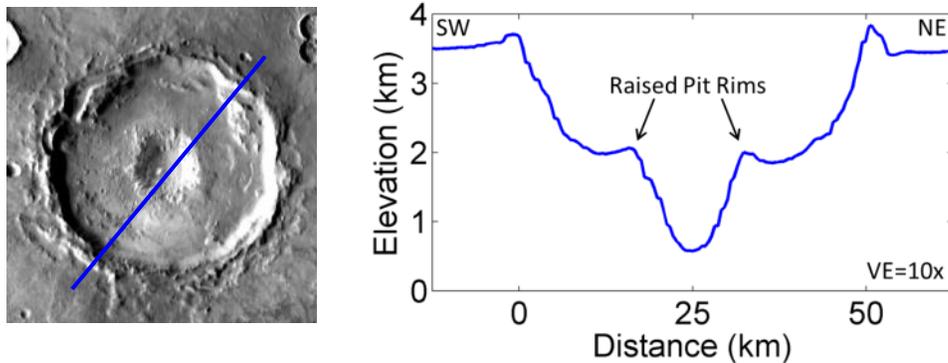
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## 41 **Background:**

42         Central pits occur in many impact structures on Mars and exhibit a crater-in-crater  
43 configuration [*e.g.*: Smith, 1976; Hodges, 1978; Barlow *et al.*, 2000; Barlow, 2010] (Fig. 1).  
44 Kilometer-scale central pits have been identified on the floors or on top of the central peaks of  
45 over 1,000 Martian impact craters with diameters as large as 125 km in diameter and down to as

46 small as 5 km in diameter [Smith, 1976; Barlow and Bradley, 1990; Barlow *et al.*, 2000; Barlow,  
 47 2011]. In our study, we focus on “floor pits” that are deeper than the surrounding floor of their  
 48 parent craters, as opposed to “summit pits” that occur atop the central peaks. Based an ongoing  
 49 survey by Barlow [2010, 2011] and this study, central floor pits have a median diameter of 0.16-  
 50 0.175 parent crater radii, such that a 50 km diameter crater might have a central pit ~8 km wide.  
 51 Their depths range from very shallow to over 1.5 km below the surrounding impact crater floor.

52



53

54 Fig. 1: THEMIS daytime IR mosaic of a 50 km diameter unnamed Martian impact crater  
 55 containing a central pit at 296.4°E, 17.6°S. A MOLA topographic profile across the center shows  
 56 typical pit morphology.

57

58 Central pit craters on Mars are confined to low and mid-latitudes, within  $\pm 70^\circ$  of the  
 59 Martian equator [Hodges *et al.*, 1980; Barlow, 2011; Garner and Barlow, 2012]. Central pits are  
 60 also common for impact craters on icy satellites, including Ganymede and Callisto [Smith *et al.*,  
 61 1979]. Central pits are seldom observed on rocky planets other than Mars, although a few dozen  
 62 are present on Mercury [Schultz, 1988; Xiao and Komatsu, 2013] and the Moon [Croft, 1981;  
 63 Schultz, 1976a, 1976b, 1988; Xiao *et al.*, 2014].

64           The presence of water-ice is believed to be involved in typical pit formation [Hodges *et*  
65 *al.*, 1980; Croft, 1981]. Although water-ice is not stable on low-latitude Martian surfaces today  
66 [Clifford and Hillel, 1983; Mellon *et al.*, 1997; Head *et al.*, 2003], water was (and might still be)  
67 present within the upper few kilometers of the surface even at low latitudes earlier in Mars'  
68 history. The possibility of significant subsurface water in pre-impact terrains is supported by the  
69 presence of layered ejecta surrounding many fresh Martian impact craters [Carr *et al.*, 1977;  
70 Gault and Greeley, 1978; Wohletz and Sheridan, 1983; Barlow *et al.*, 2000; Baloga *et al.*, 2005]  
71 and Mars Odyssey Gamma Ray Spectrometer spectra [Boynton *et al.*, 2007]. However, the  
72 process(es) responsible for forming central pits in impact craters and the role of water are still  
73 debated, and several mechanisms for pit formation have previously been proposed.

74           Wood *et al.* [1978] proposed explosive decompression may volatilize a subsurface water-  
75 rich layer, causing steam explosions and removing the core of central peaks. However, this  
76 model suffers from the difficulty of keeping water vapor from escaping early in the impact  
77 process before a central pit can be preserved [Croft, 1981; Pierazzo *et al.*, 2005; Senft and  
78 Stewart, 2011; Elder *et al.*, 2012].

79           Croft [1981], Bray [2009], Senft and Stewart [2011], Alzate and Barlow [2011] and Elder  
80 *et al.* [2012] proposed central pits could form by the melting then gravitational drainage of target  
81 water-ice through fractures underlying central uplifts. However, raised rims are also associated  
82 with many Martian central pits [Wood *et al.*, 1978; Garner and Barlow, 2012] and would not be  
83 expected with collapse structures. These models also require large volumes of water to be  
84 drained, which is unrealistic for forming the central pits on the Moon and Mercury.

85           Passey and Shoemaker [1982], Bray *et al.* [2012], and Greeley *et al.* [1982] proposed  
86 central peaks of impacts in weak target materials may collapse to form central pits. However, the

87 abundance of impact craters with central peaks and summit pits in the same regions as impact  
88 craters with floor pits suggests the target material should be strong enough to prevent collapse  
89 [Barlow, 2011].

90 Greeley *et al.* [1982] proposed and demonstrated in laboratory experiments that small-  
91 scale central pits can be excavated from impacts into layered targets causing central peak  
92 detachment. Schultz [1988] also proposed central pits are excavated as a primary result of  
93 impacts with low-velocity bolides. However, scaling up to planetary impact craters with  
94 diameters of tens of kilometers is problematic because the material-strength crater is larger than  
95 the gravity-controlled transient crater, greatly reducing the influence of any strength differences  
96 on the final crater morphology [Croft, 1981].

97 For this study, we group the previously proposed mechanisms for pit formation into those  
98 that explosively remove material upward and outward [*e.g.*: Wood *et al.*, 1978; Greeley *et al.*,  
99 1982; Schultz, 1988] and those that remove material downward [*e.g.*: Croft, 1981; Passey and  
100 Shoemaker, 1982]. During a crater-forming explosion, rocks and boulders are ejected out of the  
101 crater, layers are proximally uplifted and overturned, and ejecta is draped over the surrounding  
102 surface to create raised rims [*e.g.* Melosh, 1989]. The average grain size for ejecta decreases with  
103 radial distance from the crater, such that the largest clasts or blocks are proximal to the crater rim  
104 [*e.g.*: Gault *et al.*, 1963; O'Keefe and Ahrens, 1985; Melosh, 1989; Buhl, 2014]. Features such as  
105 sinkholes, which are typical of karst landscapes, and lava tube skylights form by gravitational  
106 collapse and do not create raised rims nor emplace material atop their rims [*e.g.*, Okubo and  
107 Martel, 1998; Salvati and Sasowsky, 2002; Cushing *et al.*, 2007; Robinson *et al.*, 2012]. The  
108 presence or absence of an ejecta blanket around central pits provides one way to distinguish

109 between explosive versus collapse scenarios for the formation of central pits. A property of  
110 ejecta blankets is decreasing average grain size with distance from an explosively-derived crater.

111

## 112 **Data and Methods:**

113 For this study, we surveyed and identified impact craters  $> \sim 10$  km in diameter  
114 containing central floor pits within  $\pm 60^\circ$  latitude of the Martian equator using the Java-based  
115 planetary geographic information system program JMARS [Christensen *et al.*, 2009]. Central  
116 pits were identified as distinctive circular depressions in the center of an impact crater that  
117 appeared to be deeper than the parent crater floor based on the available imaging and  
118 topography. Many small impact craters with diameters  $< 10$  km containing central depressions  
119 were excluded from our survey due to poor spatial resolution, as well as craters we could not  
120 confidently determine had depressions deeper than the parent floor. We excluded summit pits  
121 that occur atop central peaks and are not deeper than the parent crater floor to avoid potential  
122 bias from coherent rock or boulders on the sides of the central peaks. We also excluded  
123 ambiguous structures that might be peak rings or concentric terraces, especially in craters near  
124 the Martian simple to complex crater transition of  $\sim 6$ -7 km diameter [Garvin *et al.*, 2000, 2003].

125 Diameters were measured for both the central pits and their parent craters. Most central  
126 pits were too small to identify in the 128 pixel/deg (460 m/px) Mars Orbiter Laser Altimeter  
127 (MOLA) global mosaic [Smith *et al.*, 2001], so the  $\sim 100$  m/pixel Mars Odyssey mission Thermal  
128 Emission Imaging Spectrometer (THEMIS) [Christensen *et al.*, 2004] calibrated daytime infrared  
129 (IR) global mosaic [Edwards *et al.*, 2011] was used instead, which provides nearly complete  
130 coverage to  $\pm 60^\circ$  latitude. THEMIS daytime IR images show topography as shaded relief, since  
131 sun-facing slopes are typically warmest. Higher resolution visible images were also used to

132 observe finer-scale morphology and distinguish central morphologies that appeared ambiguous  
133 in THEMIS daytime IR. Primarily, we used Mars Reconnaissance Orbiter mission Context  
134 Camera (CTX) [Malin *et al.*, 2007; Bell *et al.*, 2013] images at  $\sim 6$  m/pixel that were map-  
135 projected and photometrically stretched from Planetary Data System (PDS) raw electronic data  
136 records, and where available we used High Resolution Imaging Science Experiment (HiRISE)  
137 [McEwen *et al.*, 2007] images at  $\sim 0.25$  to  $1.3$  m/pixel that were map-projected and  
138 photometrically stretched from PDS calibrated reduced data records. The global dust  
139 environment for central pit crater context is shown using Thermal Emission Spectrometer (TES)  
140 solar energy reflectivity (albedo) integrated from  $0.3$  to  $2.9$   $\mu\text{m}$  [Christensen *et al.*, 2001].

141         During the formation of impact and other explosive craters, coarse debris are typically  
142 ejected and scattered outside the crater. Large blocks and coarse grains have a higher thermal  
143 inertia than finer-grained materials and hold on to their heat longer through the night. This  
144 thermal inertia can be calculated from nighttime thermal images and used to estimate average  
145 grain size [Christensen, 1986]. We therefore used the THEMIS thermal inertia global mosaic as a  
146 quantitative proxy for average grain size, such that coarse-grained or blocky materials have  
147 relatively higher thermal inertias (warmer at night) while dust, sand, and other fine-grained  
148 materials have lower thermal inertias (cooler at night) [Christensen, 1986; Ferguson *et al.*, 2006;  
149 Edwards *et al.*, 2009; Edwards *et al.*, 2011]. THEMIS nighttime images and thermal inertias  
150 have previously been used to identify blocky ejecta rays from impact craters on Mars that  
151 otherwise show little or no albedo variation in visible images but where grain size trends are seen  
152 with respect to distance from the crater [McEwen *et al.*, 2005; Tornabene *et al.*, 2006]. Central  
153 pits with an annulus or a geographically skewed patch of higher thermal inertia material nearer

154 the pit rim than more distally across the surrounding parent crater floor may be classified as  
155 having a fining average grain size with radial distance, consistent with ejecta.

156 To measure the trend of thermal inertias, we circumferentially averaged the THEMIS  
157 thermal inertia mosaic over central pit craters in intervals of 0.1 parent crater radii. Because most  
158 central pits are <0.2 crater radii, we compare pit-proximal averaged thermal inertia values within  
159 the interval from 0.2-0.3 crater radii versus more distal averaged thermal inertia values at 0.5-0.6  
160 crater radii. A Student's t-test is then performed on the differences between proximal and distal  
161 averaged thermal inertias for the population of central pits. A significance level of  $P \geq 0.05$  would  
162 be deemed not statistically significant and serves as our null hypothesis: thermal inertia and  
163 average grain size do not decrease radially away from pit rims. For  $P < 0.05$ , a radial decrease in  
164 thermal inertia with distance from the pit rim would be deemed statistically significant and we  
165 would reject the null hypothesis and support an alternative hypothesis that ejecta surrounds  
166 central pits.

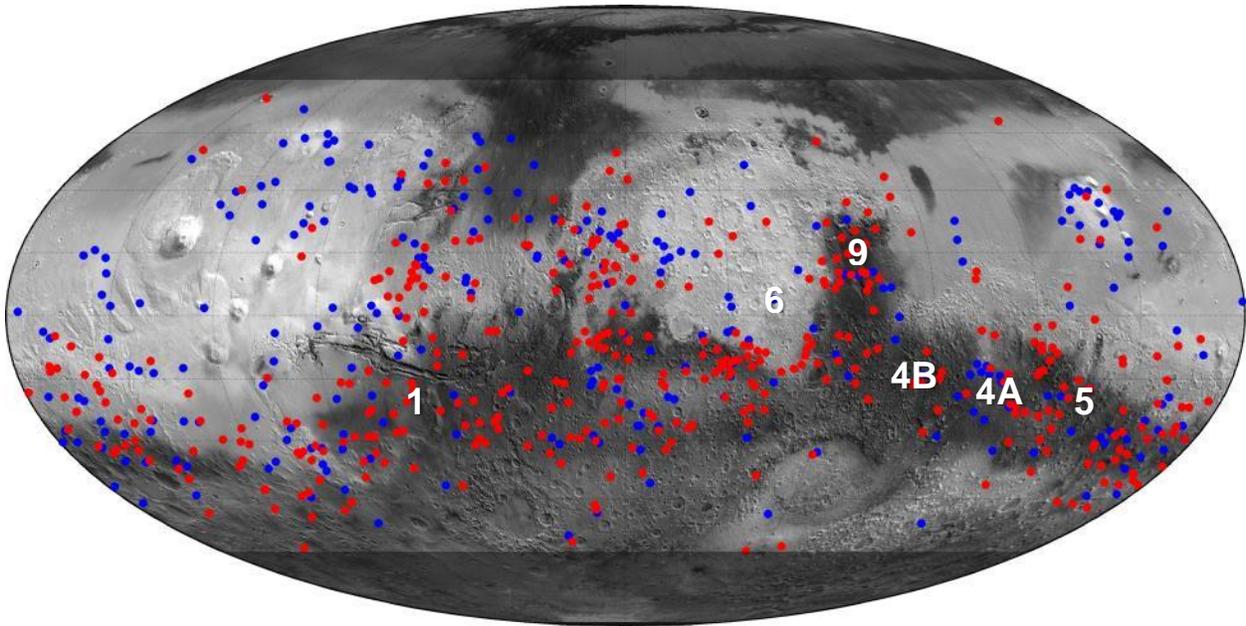
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## 168 **Results:**

169 We identified central floor pits within 654 parent craters ~10 km diameter or larger  
170 between  $\pm 60^\circ$  latitude of the Martian equator (Fig. 2). Additional smaller craters with central pits  
171 exist [Barlow, 2010, 2011], but are not well-resolved in the THEMIS thermal images used for  
172 this study. MOLA topographic profiles have very coarse resolution and may only provide insight  
173 to the largest central pit craters (Fig. 1), although complete and partially rimmed pits frequently  
174 occur in the highlands terrains [Garner and Barlow, 2012]. We identified central pits in parent  
175 impact craters with diameters ranging from ~8 to 114 km, with 95% of those parent craters being  
176 <50 km in diameter and excluding smaller potential central pit craters. The surveyed central pits

177 have a median diameter ratio to their parent craters of 0.175 with a standard deviation of 0.037  
178 (Fig. 3). These results are comparable to the median ratio of 0.16 found by Barlow [2011].

179



180

181

- Pits w/ Warm Material (n=395, 62%)
- Pits w/o Warm Material (n=240, 38%)

182

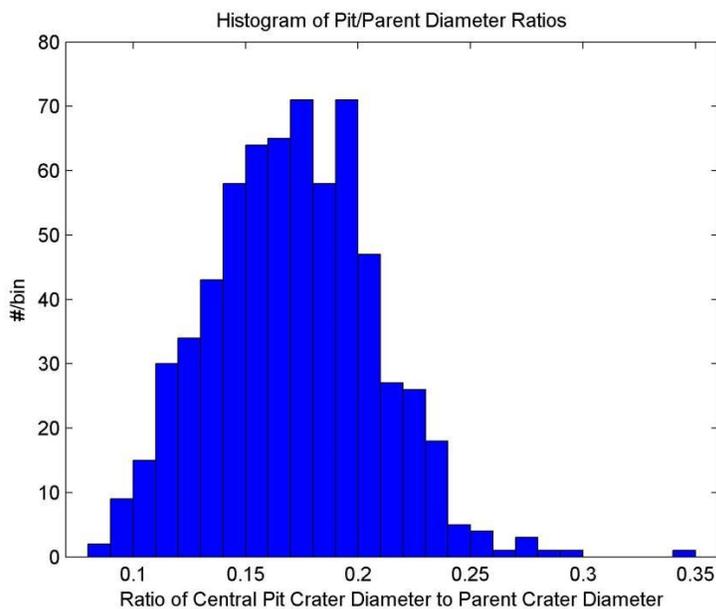
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Fig. 2: Distribution of 654 central pit craters identified in our survey of the THEMIS daytime global mosaic, within  $\pm 60^\circ$  degrees of the Martian equator, overlain on the TES albedo basemap [Christensen *et al.*, 2001] and presented in a Mollweide equal area projection. Locations of Figs. 1, 4A, 4B, 5, 6, and 9 are highlighted.



187

188 Fig. 3: Histogram showing the range of diameter ratios between central pits and their parent

189 craters that we measured. The median value is 0.175 with a standard deviation of 0.037.

190

191 Based on THEMIS-derived thermal inertias, most central pits showed higher thermal

192 inertia (coarser) material near their rim than more distally on the parent crater floor (*e.g.* Fig. 4).

193 635 of the 654 central pits had thermal images over their parent crater floors, of which 395

194 (62%) had higher average proximal thermal inertias outside the pits (between 0.2-0.3 crater radii)

195 than more distally (between 0.5-0.6 crater radii), and 240 (38%) had the same or lower thermal

196 inertia proximally than distally. Notably, 76% (254 out of 333) of central pits with diameters &gt;20

197 km have radially decreasing thermal inertia trends, 80% (175 out of 216) of central pits with

198 near-pit thermal inertia values &gt;300 TIU (less dusty) have radially decreasing thermal inertia

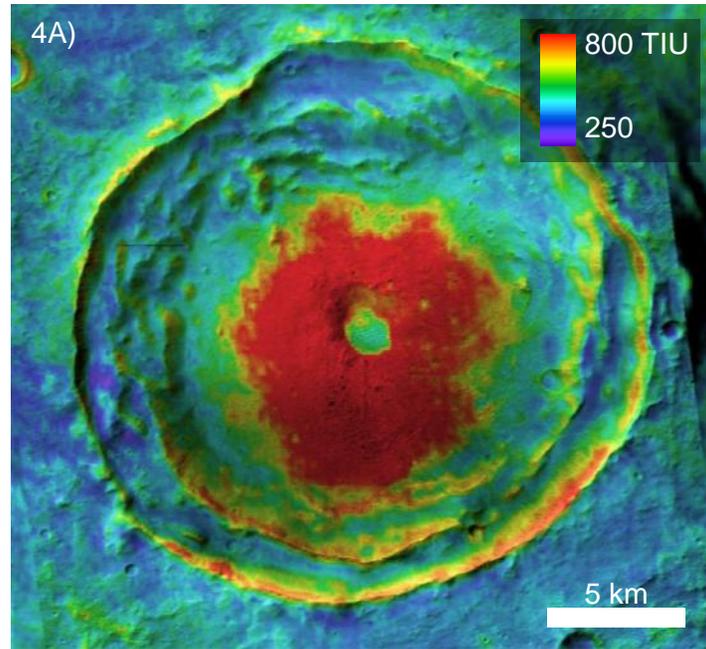
199 trends, and 89% (74 out of 83) of central pits satisfying both of the above selection criteria have

200 radially decreasing thermal inertia trends. Pits with proximal high and radially decreasing

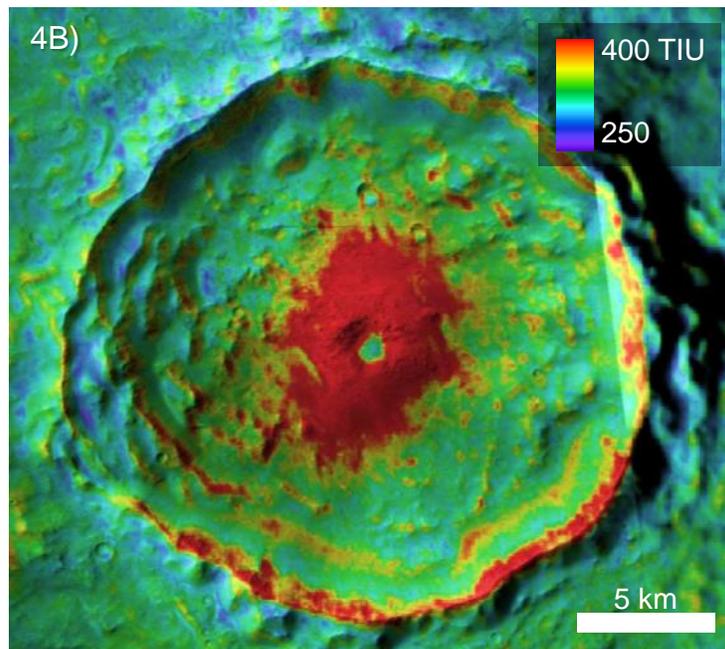
201 thermal inertias in THEMIS images sometimes show large blocky debris (up to tens of meters

202 wide) in visible CTX and HiRISE images (Fig. 5), while pits that did not show proximally high  
203 nor decreasing thermal inertias typically appeared blanketed or mantled (Fig. 6).

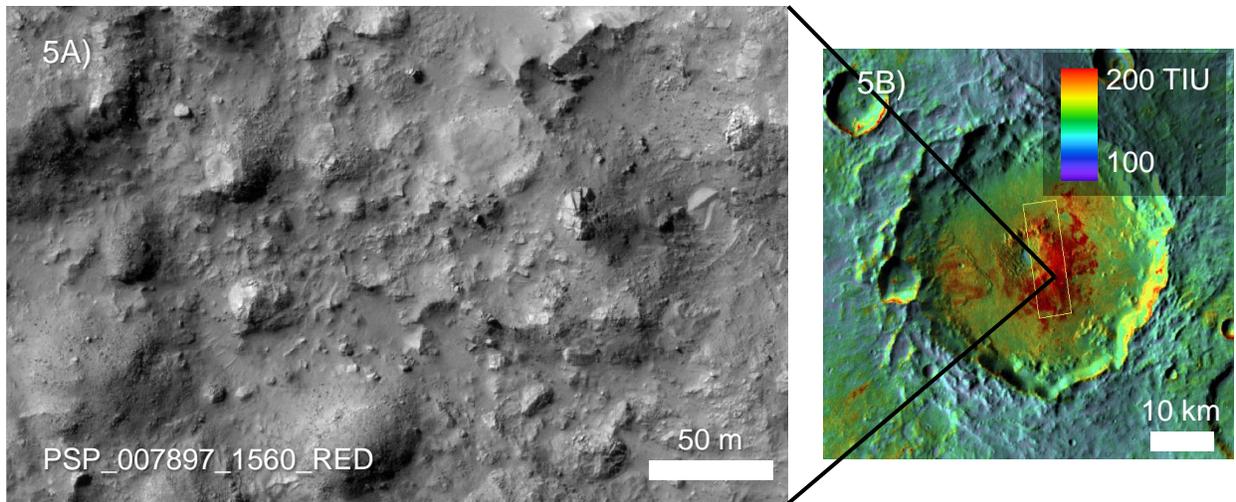
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205



206 Fig. 4: THEMIS nighttime (color) and CTX visible (shading) images showing radially  
207 decreasing high thermal inertia material interpreted as ejecta surrounding two central pit craters  
208 at A) 18.4°S, 102.7°E, and B) 14.9°S, 93.2°E. Color scales indicate thermal inertia values.



209

210 Fig. 5: A) HiRISE image showing large blocks near a central pit crater at 23.8°S, 126.8°E. B)

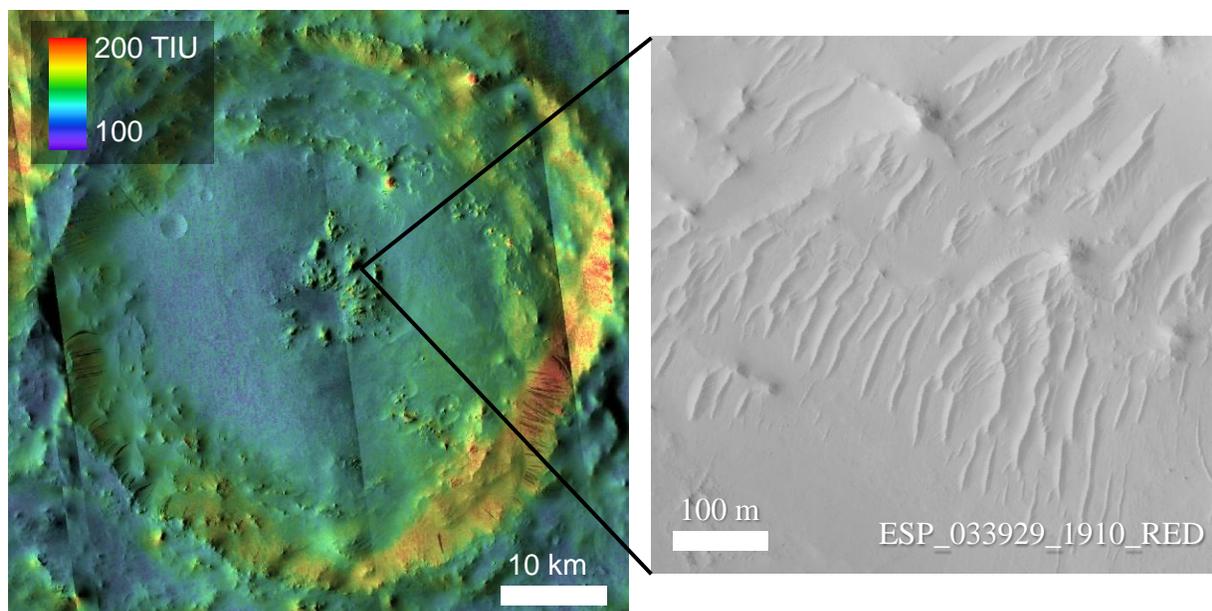
211 THEMIS nighttime IR (color) over daytime IR (shading) context image showing high-thermal

212 inertia material inferred as being blocky and confirmed by the HiRISE image. Black lines

213 indicate location of A. Yellow box in B indicates footprint of HiRISE image.

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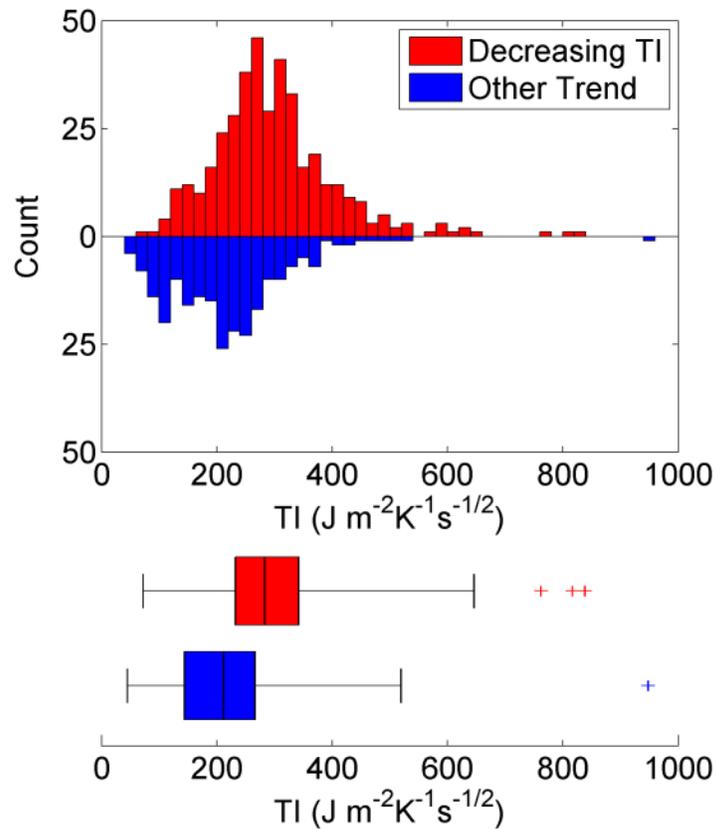


216  
217 Fig. 6: A) THEMIS nighttime IR (color) over CTX visible (shading) image showing a central pit  
218 crater at 10.9°N, 50.8°E without a radially decreasing thermal inertia. Average thermal inertia  
219 values are uniformly low across the crater floor and associated with a coating of fine-grained  
220 dust. B) HiRISE visible image enlargement of an area near the central pit showing low-contrast  
221 dust mantling the terrain.

222  
223 We conducted a paired Student's t-test to determine the confidence interval of the  
224 measured thermal inertia decreases from 0.2-0.3 crater radii to 0.5-0.6 crater radii. For the 635  
225 central pit craters with thermal images, the t-test returns a  $P < 0.01$  indicating extreme statistical  
226 significance. We therefore reject our null hypothesis that thermal inertia and average grain size  
227 do not decrease radially away from pit rims, and support an alternative hypothesis that pits are  
228 surrounded by ejecta with grain size decreasing with distance away from the pit.

229 Central pits in Tharsis, Elysium, Arabia, and other dusty regions, characterized by high  
230 TES albedos and low thermal inertia values, tend to not be surrounded by material with radially  
231 decreasing thermal inertia trends (Fig. 1). The median proximal thermal inertia for central pits

232 with radially decreasing thermal inertias is 283 thermal inertia units ( $1 \text{ TIU} = 1 \text{ J m}^{-2} \text{ K}^{-1} \text{ s}^{-1/2}$ ) with  
 233 a standard deviation of 121 TIU, while the median proximal thermal inertia for central pits with  
 234 other, radially non-decreasing thermal inertia trends is 205 TIU with a standard deviation of 145  
 235 TIU (Fig. 7).

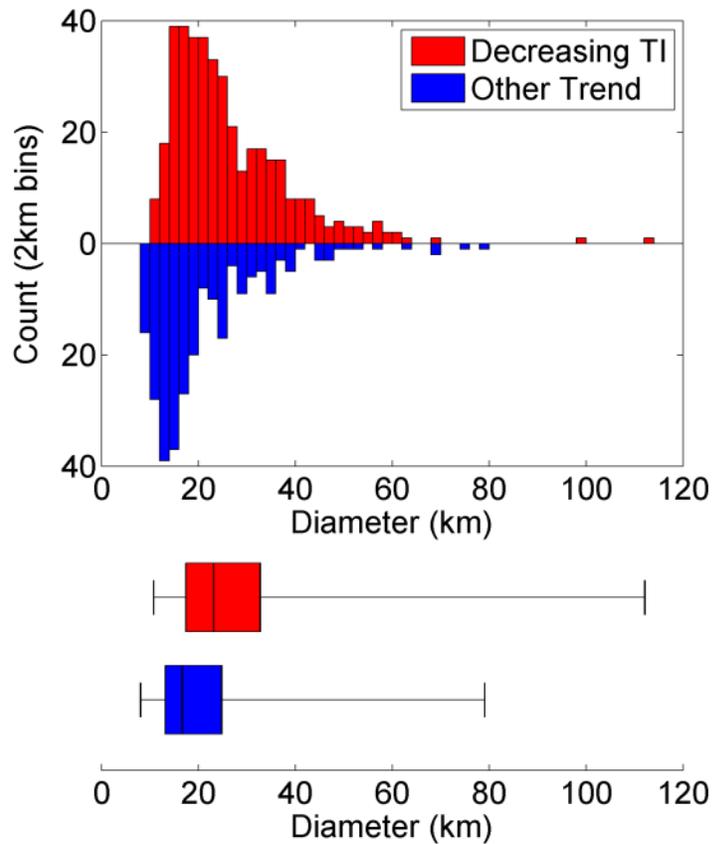


236

237 Fig. 7: Histogram and box-and-whisker plot of central pit craters exhibiting radially decreasing  
 238 thermal inertia trends (red) and radially non-decreasing thermal inertia trends (blue) plotted  
 239 against THEMIS thermal inertia values. Lower thermal inertias are indicative of finer average  
 240 grain size and dustiness.

241

242 Smaller central pits also tend not to show radially decreasing thermal inertias (Fig. 8).  
 243 Based on the population of impact craters observed with THEMIS data, the median diameter for  
 244 parent craters containing pits with warm material is  $\sim 23.3$  km and the median diameter for  
 245 craters with pits lacking it is  $\sim 16.7$  km, both cases being above the simple/complex transition of  
 246 6-7 km for Martian craters [Garvin *et al.*, 2000, 2003].



247  
 248 Fig. 8: Histogram and box-and-whisker plot of craters containing central pits exhibiting radially  
 249 decreasing thermal inertia trends (red) and radially non-decreasing thermal inertia trends (blue)  
 250 plotted against parent crater diameter.

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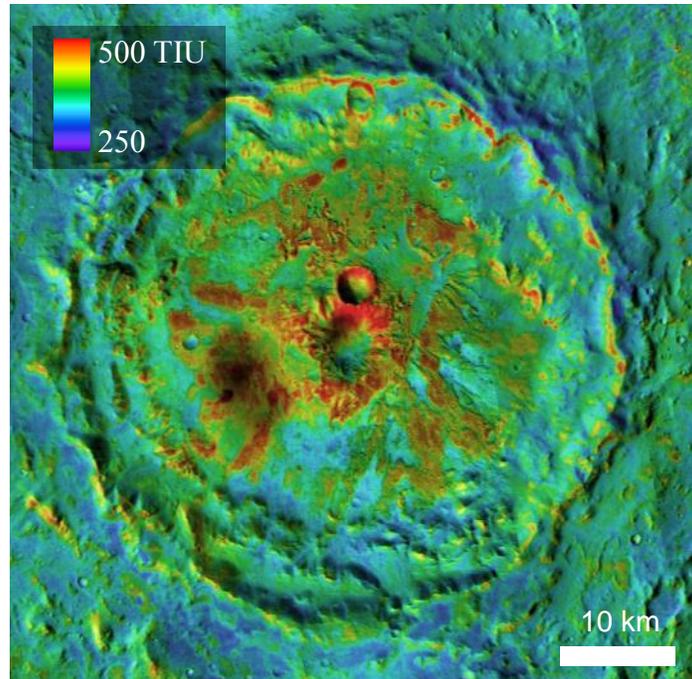
254 **Discussion:**

255           The raised rims around some pits [Wood, 1978; Garner and Barlow, 2012] are suggestive  
256 of explosive excavation, similar to their parent craters, which also have raised rims. As discussed  
257 by Garner and Barlow [2012], raised rims are more frequently observed in larger central pits than  
258 smaller ones. They also argue that the preferred distribution of rimmed pits in highlands regions  
259 and non-rimmed pits in volcanic plains suggests that target material strength and/or volatile  
260 content may also limit the expression of raised rims. Some very small scale pits on Mars  
261 believed to have formed from volatile release in impact melt have been identified and also  
262 exhibit slightly raised rims, although they are not exclusive to crater centers and do not exhibit  
263 well-defined ejecta blankets [Tornabene *et al.*, 2012; Boyce *et al.*, 2012]. Surfaces visible in  
264 some CTX and HiRISE images show large (meter-scale) blocks in warm patches adjacent to  
265 central pits (*e.g.*, Fig. 5), consistent with the expected correlation between warm material and  
266 coarse surfaces. Such blocks and megablocks are commonly observed near explosively-formed  
267 craters, including at the Ries crater in Germany [*e.g.* Gault *et al.*, 1963], as well as at some  
268 Martian craters [*e.g.*, Caudill *et al.*, 2012]. Combined with the spatial correlation of warm  
269 material and central pits, we interpret the blocks scattered around central pits to be explosively-  
270 emplaced pit ejecta.

271           The observability of high thermal inertia, coarse-grained material appears linked to the  
272 size of the pit. Small craters excavate smaller volumes of material that is finer-grained on  
273 average than larger craters [*e.g.*: Gault *et al.*, 1963; O’Keefe and Ahrens, 1985; Melosh, 1989;  
274 Buhl, 2014]. Fine-grained rocks are more easily eroded or buried than coarser-grained rocks, so  
275 the coarser ejecta at larger pits should be preferentially preserved and less buried. Surface diurnal  
276 thermal inertias are sensitive to materials within a few thermal skin depths (several centimeters)

277 of the surface, so any ejecta would have to be buried by no more than a few centimeters of dust  
278 in order to be observable. Accumulated dust and sand is frequently observed on Mars and is  
279 indicated in our analysis as low thermal inertia values due to dust's fine grain size (Fig. 6). The  
280 smaller grain size distribution of ejecta for smaller craters is therefore expected to decrease the  
281 positive detection of ejecta using diurnal thermal inertias.

282         The presence of high thermal inertia material on parent crater floors near pits would not  
283 necessarily need to be due to ejecta. To avoid many false-positives, we have calculated the trend  
284 in thermal inertia (grain size) with radial distance from the pit. For example, post-impact lava or  
285 perhaps impact melt flows occur on the floors of some craters containing central and have high  
286 thermal inertias, although small flow lobes are easily distinguishable (Fig. 9), and more extensive  
287 lava or impact melt flows could potentially fill central pits. We expect impact melt ponds to be  
288 distributed throughout the crater floor, so measuring a radially decreasing trend in thermal inertia  
289 as opposed to only using high thermal inertia values avoids this problem in most cases. Patchy or  
290 partial erosional uncovering of consolidated parent crater fill rocks could also explain higher  
291 thermal inertias relative to the surrounding crater floor; however, we consider the selective  
292 removal of significant amounts of dust from the centers of parent craters, but not in the dusty  
293 plains surrounding many parent craters, to be unlikely. Additionally, significant erosion on the  
294 parent crater floor is inconsistent with the presence and preservation of raised rims around many  
295 central pits. Thermal inertias are also low for relatively fine-grained aeolian dunes or other  
296 bedforms that often form in the centers of craters, and confirmed in CTX and HiRISE images  
297 (Fig. 6).



298

299 Fig. 9: THEMIS nighttime IR (color) over CTX visible (shading) image showing high thermal  
300 inertia lava flow lobes (red, oranges, and yellow irregular bands on crater floor) on the floor of  
301 an impact crater containing a central pit at 28.5°N, 83.4°E.

302

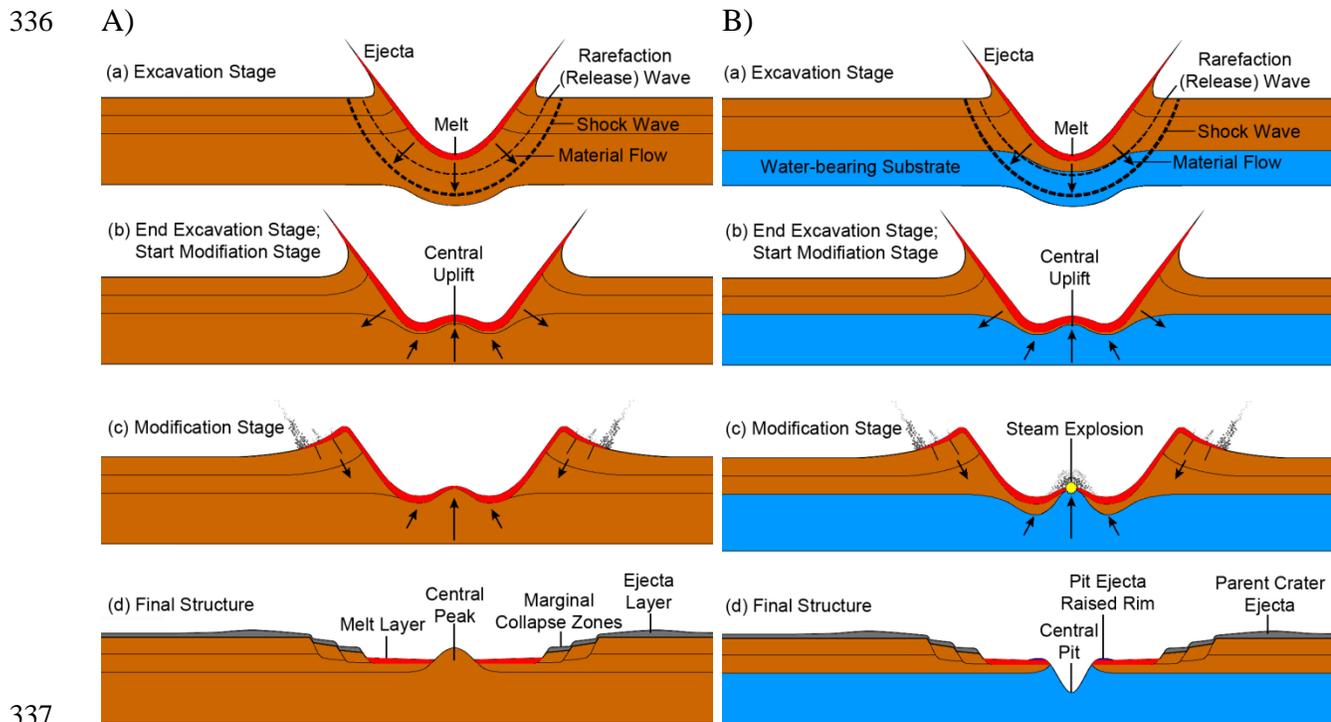
303 The results of our thermal inertia study are consistent with and support both the Wood *et*  
304 *al.* [1978] and Greeley *et al.* [1982] models; however, each suffers from a critical weakness. The  
305 Wood *et al.* [1978] model for an explosive pit origin suffers from the difficulty of keeping vapor  
306 from escaping early in the impact process before a pit can be preserved. The Greeley *et al.*  
307 [1982] central peak detachment model also suffers from issues scaling up from the laboratory to  
308 planetary impact craters. Alternatively, an explosive reaction could potentially result from  
309 mixing of water-ice and molten rock through several mechanisms. For example, a post-impact  
310 magmatic intrusion could intrude into a crater and react with the ground water as a maar volcano  
311 [Wohletz, 1986; Begét *et al.*, 1996]; however, we would not expect such a scenario to  
312 consistently form pits in crater centers. Heavy fracturing and brecciation during the impact

313 process may allow fluids (either impact melt, or liquid water) to mobilize and permeate the  
314 substrate and come into contact with each other, similar to the fluid flow described by Elder *et*  
315 *al.* [2012]. Although liquid water may move freely through fractures, Elder *et al.* finds that  
316 impact melt would cool too quickly due to its high melting temperature and larger temperature  
317 difference with the country rock. Rain or ice-bearing fallback ejecta could also be deposited on  
318 top of impact melt pools or suevite deposits [Segura *et al.*, 2002], but that would not necessarily  
319 require that pits always form in the centers of their parent craters, nor that they be consistently  
320 sized. Below, we describe an alternate model for bringing water into contact with impact melt.

321

### 322 **Alternate Pit Formation Model:**

323 We present an alternate hypothesis that -- unique among other pit origin hypotheses --  
324 predicts an explosion late enough in the impact process for central pits to be preserved and has a  
325 properly scaled analog. In our uplift contact model, impact central uplifts bring water (as liquid,  
326 ice, or both) vertically up and into contact with near-surface impact melt to initiate late-stage  
327 steam explosions and form central pits (Fig. 10). Central uplift occurs late in the impact process  
328 during the modification stage, after most crater fill has settled [*e.g.*, Melosh, 1989]; thus, pit  
329 formation concurrent with central uplift is consistent with the apparent lack of infilling of deep  
330 pits. As we describe in the next paragraph, our explosive central pit model is akin to an inverted  
331 maar volcano [*e.g.* White and Ross, 2011], except instead of magma rising up into contact with  
332 groundwater or permafrost, a water-bearing substrate is uplifted into contact with impact melt.  
333 Similarly-scaled events have been observed at monogenetic maar volcanoes with diameters of up  
334 to 8 km on the Seward Peninsula in Alaska [Begét *et al.*, 1996], where the permafrost buffers the  
335 water-magma interaction to achieve high heat transfer efficiencies [Wohletz, 1986].



338 Fig. 10: Schematic cartoons illustrating steps in complex crater formation resulting in: A) a  
 339 classical central peak [modified from French, 1998], and B) our proposed new "uplift contact  
 340 model" for Martian central pit crater formation.

341

342 As the central uplift rises, it brings deeply-sourced water-bearing rock from below the  
 343 transient cavity up into contact with shallow crater fill deposits and impact melts. We would not  
 344 expect significant vertical mixing of sub-transient cavity material outside the central uplift, so  
 345 these large pits should always be in the centers of their parent impact craters. As the water-  
 346 bearing central uplift rises into contact with impact melt and other hot debris, the thermal energy  
 347 from the melt may be transferred to the water, resulting in a steam explosion to eject material  
 348 outward, raise rims, and deposit ejecta surrounding the pits (with average grain sizes decreasing  
 349 with radial distance, as we found in this study). As material is ejected outwards, the walls may  
 350 become unstable and slump hot debris and impact melt into the pit cavity. There, the new rush of

351 melt and hot rocks may again react with uplifting water to recharge the system and iteratively  
 352 trigger a series of explosions to further deepen and widen the central pit. When central uplift  
 353 slows, the vertical mixing of water decreases and the explosions will cease.

354 We explored the theoretical plausibility of whether enough thermal energy could have  
 355 been available in a post-impact environment to initiate steam explosions capable of creating  
 356 kilometer-scale central pits. We started with the empirical model shown below which predicts  
 357 the mass ratio of melted ( $m_m$ ) to displaced ( $m_d$ ) impact target materials in a silicate target (Eq. 1)  
 358 [O’Keefe and Ahrens, 1982; Melosh, 1989]:

$$359 \quad m_m/m_d = 1.6 \times 10^{-7} \times (g \times D_i)^{0.83} \times v_i^{0.33} \quad (1),$$

360 where  $g$  is planetary gravity,  $D_i$  is parent crater diameter and  $v_i$  is bolide velocity. We assign the  
 361 following values for our calculations: gravity  $g = 3.711 \text{ m/s}^2$  and bolide velocity  $v_i = 10 \text{ km/s}$   
 362 [Ivanov *et al.*, 2002]. We also assumed that any melt generated remained within the parent  
 363 crater. Finally, we modeled the parent crater as a half-ellipsoid and applied the mass fraction to  
 364 determine the volume and mass of melt produced (Eqs. 2,3):

$$365 \quad V_m = (m_m/m_d) \times (2/3) \times \pi \times d_i \times (D_i/2)^2 \quad (2) \text{ and}$$

$$366 \quad m_m = \rho_m / V_m \quad (3),$$

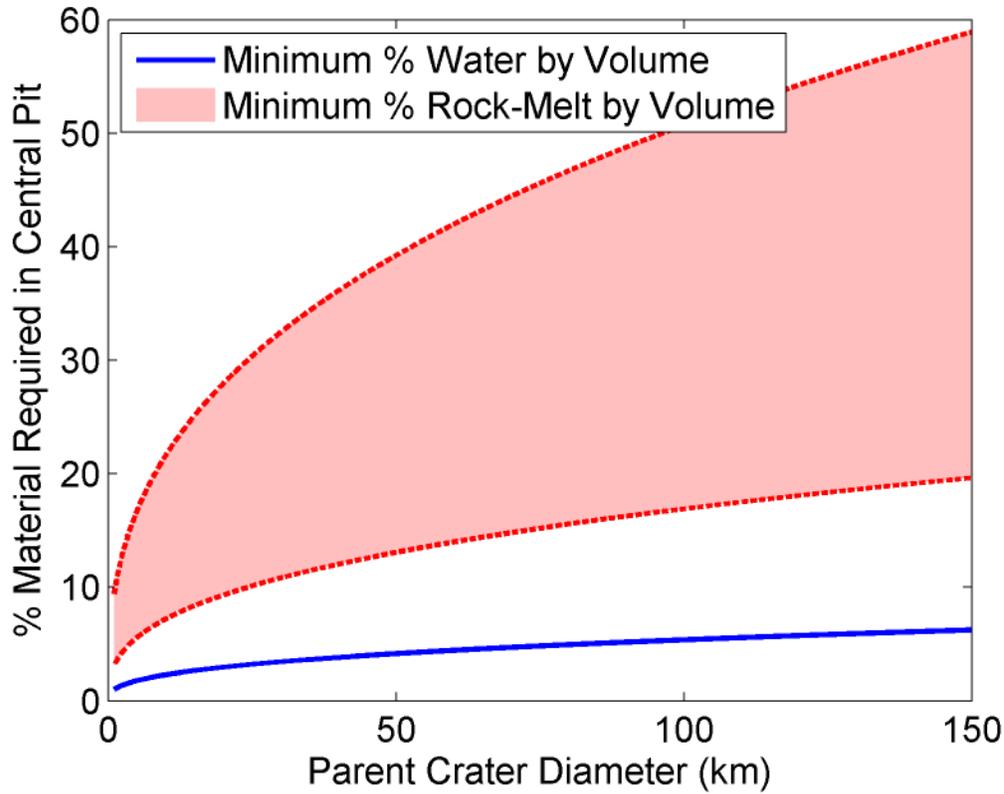
367 where  $V_m$  is the volume of melt,  $d_i$  is the depth of the parent crater, and  $\rho_m$  is the density of the  
 368 melt. We assume a depth of complex craters (in km) of  $d_i = 0.357 D_i^{0.52}$  [Tornabene *et al.*, 2013].  
 369 Sato and Taniguchi [1997] used the following empirical equation to predict the energy required  
 370 to form a crater via volcanic, nuclear, and chemical explosions, independent of origin. The  
 371 equation can similarly be applied to central pits (Eq. 4):

$$372 \quad E_c = 4.45 \times 10^6 \times D_p^{3.05} \quad (4),$$

373 where  $E_c$  is the energy of pit formation and  $D_p$  is the diameter of the pit, for which we assume a  
 374 median pit-to-parent crater diameter ratio of 0.16 [Barlow, 2010, 2011]. The total thermal energy  
 375 transfer required to melt ice and boil water to steam can be calculated using specific and latent  
 376 heats (Eq. 5):

$$377 \quad H_w = m_w \times L_f + m_w \times c_{lq} \times \Delta T_w + m_w \times L_v \quad (5),$$

378 where  $H_w$  is the energy transferred to the water,  $m_w$  is the mass of water,  $L_f$  is the latent heat of  
 379 fusion,  $c_{lq}$  is the specific heat of liquid water,  $\Delta T_w$  is the temperature change of liquid water, and  
 380  $L_v$  is the latent heat of vaporization. We assign values for  $L_f = 3.34 \times 10^5$  J/kg,  $c_{lq} = 4.187 \times 10^3$   
 381 J/kg·K, and  $L_v = 2.257 \times 10^6$  J/kg [Moran and Shapiro, 2008]. Evaluating Eq. 5, we see that an  
 382 investment of  $3.023 \times 10^6$  J is required to turn 1 kg of water from ice (273 K) to steam (373 K).  
 383 We assume that the steam is not heated to higher temperatures, although a smaller amount of  
 384 superheated steam might also satisfy the energy requirements for explosivity. The thermal  
 385 energy of vaporization, specifically the step of converting water to steam, can be transformed to  
 386 kinetic energy that can form a pit. The mass of steam required is calculated by dividing the pit  
 387 formation energy from Eq. 4 by the latent heat of vaporization. Dividing this result by the density  
 388 of ice provides the volume of ice required to form a central pit. As shown in Fig. 11, assuming a  
 389 half-ellipsoidal pit geometry with the pit depth (in km)  $d_p = 0.276 D_p^{0.68}$  [Tornabene *et al.*, 2013],  
 390 only a small amount of water (comprising 2-6% of a central pit's volume) would need to be  
 391 vaporized to form a central pit for the parent crater diameters observed (5-125 km [Barlow,  
 392 2011]).



393

394 Fig. 11: Required amounts of water and impact melt for heat energy transfer to form a kilometer-  
 395 scale (pit) crater shown as percent by volume with respect to the volume of a central pit crater.

396 The range in impact melt volume represents uncertainty due to varying heat transfer efficiency  
 397 between 0.1-0.3.

398

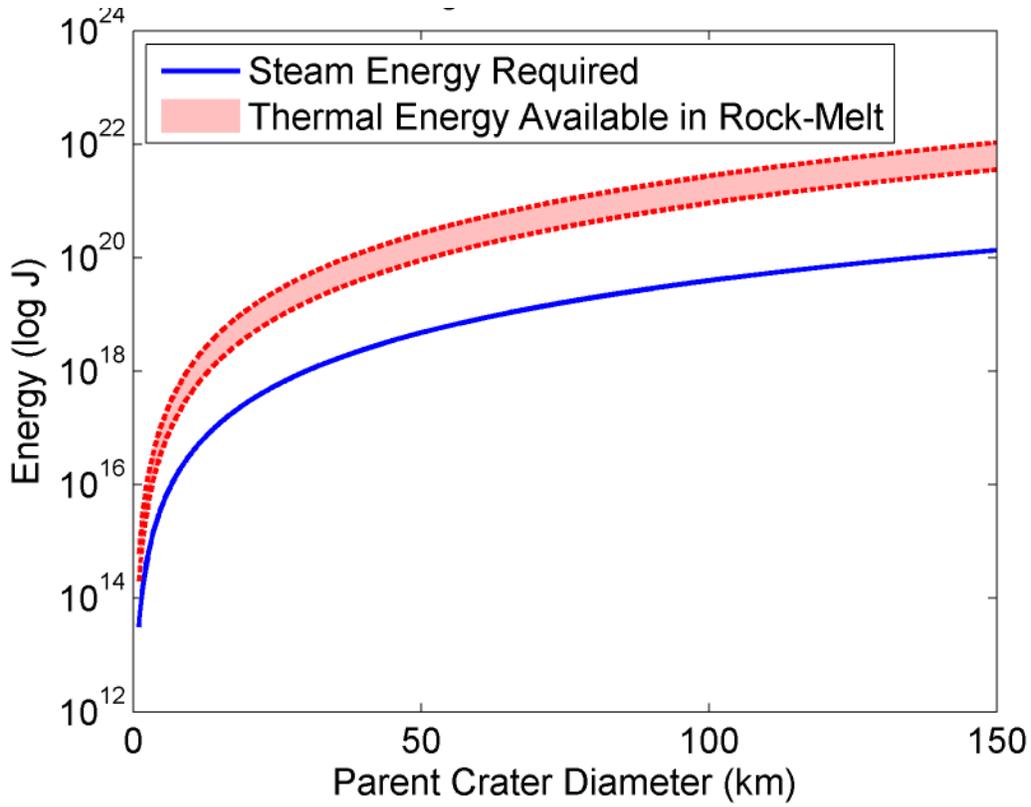
399 The amount of thermal energy available in impact melt may also be calculated using  
 400 specific heats (Eq. 6):

$$401 \quad E_m = m_m \times c_{pm} \times \Delta T_m = \rho_m \times V_m \times c_{pm} \times \Delta T_m \quad (6),$$

402 where  $E_m$  is the energy required for cooling rock,  $c_{pm}$  is the specific heat of rock,  $\Delta T_m$  is the  
 403 temperature change of the rock. We assume a basaltic melt composition and assign values of  $\rho_m$   
 404 = 2900 kg/m<sup>3</sup> [Judd and Shakoor, 1989];  $c_{pm}$  = 1000 J/kg·K [Wohletz, 1986]; and change of  
 405 temperature (from the basalt solidus to the STP boiling point of water)  $\Delta T_m$  = 1473 K – 373 K =

406 1100 K [Wohletz, 1986]. It should be noted that impact melts can also be superheated, perhaps  
407 up to 1700°C (1973 K) [Zieg and Marsh, 2005], so our calculations may underestimate the  
408 thermal energy available by ~50%. Adiabatic heat transfer efficiency is typically ~0.1 or less due  
409 to poor mixing; however, it can reach an optimal efficiency of ~0.3 for water/melt ratios of 0.3-  
410 0.5 [Wohletz, 1986]. Such optimal efficiencies are believed to be present for maars in  
411 permafrost, as suggested by the largest, kilometer-scale terrestrial maars found in the Seward  
412 Peninsula, Alaska [Begét *et al.*, 1996]. Our calculations consider cases with both 0.1  
413 (suboptimal) and 0.3 (optimal) efficiencies.

414         The mass of impact melt required to vaporize ice to steam can be calculated by setting the  
415 total heat transfer  $H_w$  from Eq. 5 equal to the product of the heat transfer efficiency and the  
416 impact melt thermal energy from Eq. 6. As shown in Fig. 11, the impact melt must comprise a  
417 volume greater than or equal to 6-18% of the central pit's volume for an optimal thermal  
418 efficiency of 0.3, or 17-55% of the central pit's volume for a suboptimal thermal efficiency of  
419 0.1. The total energy transfer required for vaporizing ice ( $H_w$ ) from Eq. 5 can also be compared  
420 to the total energy available from impact melt by multiplying Eq. 6 with the value(s) for heat  
421 transfer efficiency (Figs. 12,13). Based on these calculations, sufficient thermal energy should be  
422 available via impact melt to vaporize small amounts of ice that act explosively to form central  
423 pits within kilometer-scale impact structures. However, not all Martian craters exhibit central  
424 pits. Below, we discuss the material requirements may inhibit the explosive formation of some  
425 central pits on Mars.



426

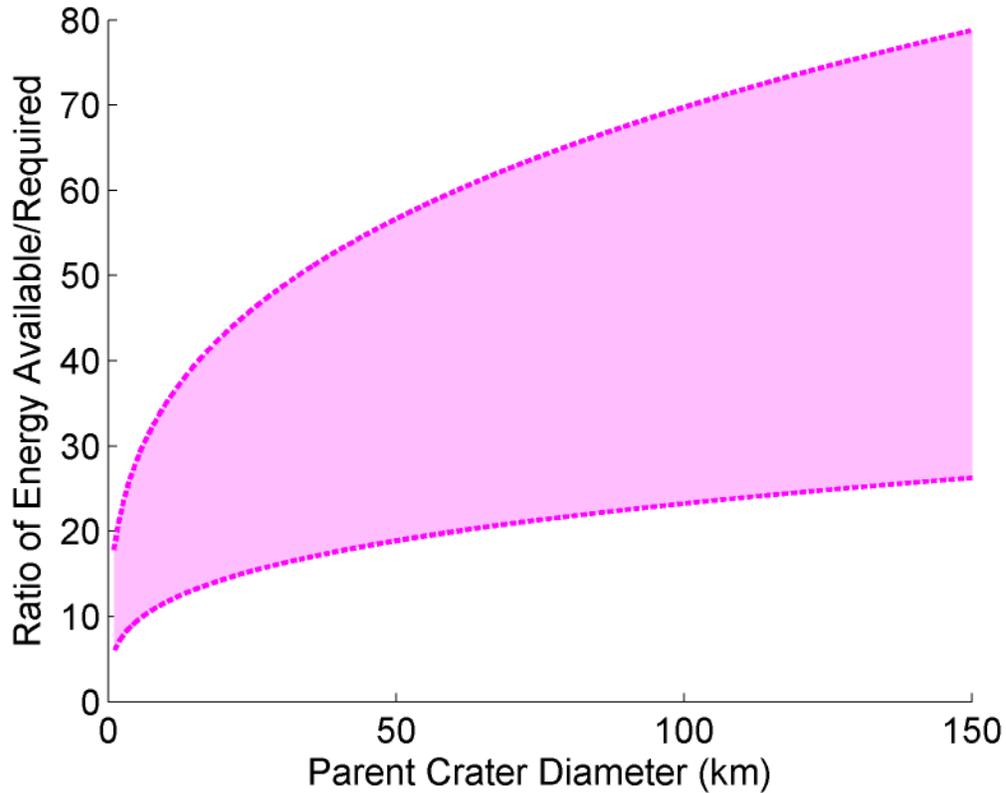
427 Fig. 12: Thermal energies of water required to convert ice to steam to provide the energy for

428 creating central pit craters (blue line) of differing diameter. Also shown is the available thermal

429 energy from impact melt, after applying thermal efficiency values of 0.1 (lower red curve) to 0.3

430 (upper red curve).

431



432

433 Fig. 13: Ratios of available / required thermal energy for vaporizing enough steam to explode  
434 and form a central pit, with respect to crater diameter. The range of in energy ratios reflects  
435 variations in heat transfer efficiency over a range of 0.1 (lower curve) to 0.3 (upper curve).

436

437 First, an appropriate volume of water must be available in the central uplift. If too little  
438 water (or too low a concentration) is present, there may not be sufficient steam to form a large  
439 pit. Even if water was initially present in the target rocks, large impacts (with crater diameters of  
440 several tens to hundreds of km) likely remove most subsurface volatiles early in the impact  
441 process such that not enough water is available to react with the impact melt to form a pit.  
442 Conversely, if the system has excess water, there may not be enough thermal energy in the  
443 impact melt to heat the excess water and still vaporize enough to sustain an explosion and make  
444 a pit.

445           Second, an appropriate volume of impact melt must be retained within the parent impact  
446 crater. Smaller impact craters produce less melt proportionally and distribute that impact melt  
447 more sparsely, so small craters may not have enough consolidated impact melt even if enough  
448 water is present. Larger impact craters might also produce excess impact melt that could fill in  
449 any central pits that might form. A similar phenomenon is thought to have occurred at the  
450 Sudbury impact structure, where steam explosions created the brecciated Onaping Formation but  
451 the explosive depressions themselves were filled in and erased [Grieve *et al.*, 2010]. Another  
452 interesting aspect of the uplift contact model is that since our calculations show it only requires  
453 small amounts of water (perhaps as little as 2-6% by volume), it provides a possible explanation  
454 for the formation of the small number of central pits observed on Mercury [Schultz, 1988; Xiao  
455 and Komatsu, 2013] and the Moon [Croft, 1981; Schultz, 1976a, 1976b, 1988; Xiao *et al.*, 2014],  
456 which should have insufficient water or other volatiles to form by drainage and collapse models  
457 [*e.g.* Croft, 1981]. Although we did not measure summit pit-related thermal inertias in our  
458 survey, summit pits would be expected to form as in our uplift contact model when steam  
459 explosions start but become water- or impact melt-limited. In such a case, the explosive reaction  
460 fails before uplift has ceased and an incomplete pit is left superposed on a remnant central peak.

461           Based on our uplift contact model, we propose the following testable predictions. First, a  
462 the ejecta deposit is expected to contain abundant fractured and fragmented glassy impact melt,  
463 similar to the Onaping Formation at Sudbury [Grieve *et al.*, 2010]. This layer of glassy deposits  
464 should overlay more coherent impact melt deposits. Second, lithic clasts and mineral  
465 assemblages found stratigraphically below the transient crater should be found on the floor of the  
466 parent crater, with the greatest abundance proximal to the rim. Third, the stratigraphic sequence

467 of rocks around central pits should be overturned. Finally, *in situ* measurements of material  
468 around the pit should show decreasing average grain sizes with radial distance from central pits.

469

470 **Conclusions:**

471 The presence of raised rims and blocky material surrounding Martian central pits are  
472 suggestive of ejecta from an explosive pit origin. Over 60% of all central pits in our global  
473 survey have material with radially decreasing thermal inertias around them, and 89% of central  
474 pits craters with diameters >20 km and non-dusty proximal thermal inertias >300 TIU have  
475 radially decreasing thermal inertias. The population of central pit craters as a whole has a  
476 statistically significant ( $P < 0.01$ ) decrease in thermal inertia radially outwards from pit rims. We  
477 interpret these findings as a typical decrease in average grain size with increasing distance away  
478 from central pits. As expected, dust masks the diurnal thermal signature around many central  
479 pits. This effect is amplified in smaller pits due to their less voluminous and finer-grained ejecta  
480 that are more easily buried or eroded. Previously proposed models do not satisfactorily explain  
481 all observed characteristics of central pits. We have therefore proposed a new "uplift contact  
482 model" to explain the observed morphologies (i.e., geometries, raised rims) and thermal  
483 properties (radially decreasing thermal inertias/average grain size) of Martian central pit craters.  
484 Our thermal calculations show that only  $\geq 2$ -6% water by volume is required to create a  
485 phreatomagmatic explosion and form central pits. Our explosive origin model is also  
486 advantageous over drainage and collapse models in explaining the small number of central pits  
487 on Mercury and the Moon using only minor amounts of volatiles in localized pre-impact  
488 subsurfaces. Drainage and collapse may still be a viable method for pit formation on icy

489 satellites, but an explosive origin appears to be the preferred mechanism on Mars (and other  
490 rocky planets) for forming central pit craters.

491

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497

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